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Математическое моделирование

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RESEARCH ARTICLE

Analytical model for the normal component of magnetic induction of a permanent magnets

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Abstract

Objectives. In a measuring system based on the inductive transmission of information from a moving structure to a stationary signal receiver, the signal carrying useful information about the parameters of the moving structure is formed by a magnetic system containing a permanent magnet mounted on the stationary part of the measuring system. The magnetic field of the permanent magnet (MFPM) determines the magnetic flux, and, consequently, the induction current in a conducting coil located on the moving structure. In order to theoretically justify the parameters of the measuring system including the optimization of its components, a simple and easy-to-use analytical model of the useful signal for determining the requirements for the mathematical description of the MFPM is required. The use of known solutions for developing an analytical model of the useful signal of the measuring system is complicated by the need to use inverse trigonometric functions or the results of numerical calculations. The present work sets out to obtain an exact solution to the problem of calculating the MFPM and on this basis to develop a simple, convenient analytical model of the normal component of the magnetic induction vector (NCMIV) of a permanent magnet used to develop an analytical model of the useful signal.

Methods. The equivalent solenoid method was used along with mathematical analysis approaches.

Results. An exact solution for calculating the normal component of the magnetic induction vector of the parallelepiped-shaped permanent magnet was obtained. Based on this, a straightforward and easy-to-use analytical model of the NCMIV was developed, which closely approximates the formula derived for the exact solution.

Conclusions. The developed analytical model of the NCMIV can be used for theoretical development of an analytical model of the useful signal of a measuring system with inductive transmission of information about the parameters of a moving structure to a stationary signal receiver.

Keywords: permanent magnet, magnetic induction, equivalent solenoid, normal component, analytical model

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НАУЧНАЯ СТАТЬЯ

Аналитическая модель нормальной составляющей магнитной индукции постоянного магнита

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Резюме

Цели. В измерительной системе с индукционной передачей информации с перемещающейся конструкции на неподвижный приемник информационный сигнал, несущий информацию о параметрах перемещающейся конструкции, формируется магнитной системой, содержащей постоянный магнит, установленный на неподвижной части измерительной системы. Магнитное поле постоянного магнита (МПМ) определяет магнитный поток, и, следовательно, индукционный ток в другом элементе магнитной системы – проводящем витке, расположенном на перемещающейся конструкции. Для теоретического обоснования параметров измерительной системы, в т.ч. для оптимизации ее составных частей, необходима простая, удобная для применения аналитическая модель информационного сигнала (АМИС), что определяет требования к математическому описанию МПМ. Известные решения задач по расчету МПМ содержат обратные тригонометрические функции или представлены результатами численных расчетов, что затрудняет их использование для разработки АМИС измерительной системы. Целью данной статьи является получение точного решения задачи расчета МПМ и разработка на основании этого точного решения аналитической модели нормальной составляющей вектора магнитной индукции (НСВМИ) постоянного магнита, используемой для разработки АМИС.

Методы. Использовались методы математического анализа и метод эквивалентного соленоида.

Результаты. Получено точное решение задачи расчета НСВМИ МПМ, имеющего форму параллелепипеда, на основании которого получено выражение, аппроксимирующее формулу точного решения, – аналитическая модель НСВМИ.

Выводы. Полученная аналитическая модель НСВМИ может быть использована для теоретической разработки АМИС измерительной системы с индукционной передачей информации о параметрах перемещающейся конструкции на неподвижный приемник сигнала.

Ключевые слова: постоянный магнит, магнитная индукция, эквивалентный соленоид, нормальная составляющая, аналитическая модель

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INTRODUCTION

A radio channel is mostly used to convey information about the parameters of a structure that changes position relative to a stationary signal receiver (consider, for instance, relaying gas pressure readings within a vehicle's tire). In this situation, measurement transducers are mounted on the structure to transform characteristics

defining its status into electrical impulses, which are subsequently transmitted to a radio transmitter. A radio signal receiver is mounted on the stationary part, such as the body of the vehicle, close to where the pneumatic tire wheel changes position.

Transmitting information via a radio channel often encounters challenges stemming from limitations imposed on the use of electromagnetic radiation. In

a measuring system with inductive transmission of information about the state of a moving structure [1, 2], the useful signal carrying information about the state of the structure is formed by a magnetic system, which is formalized as shown in Fig. 1. A rectangular conducting coil (1) is positioned on the upper edge of a permanent magnet (PM), which is configured as a parallelepiped with a height of H and the base having dimensions a and b . In a plane parallel to the PM upper edge, a structure with a placed-in rectangular coil 2 travels relative to the magnet at velocity V , with the center of coil 2 (point O_2) projected onto the Ox axis in a manner that its projection (point A) traverses along the Ox axis. The distance between the plane and the upper edge of the magnet is equal to d . The electrical resistance of coil 2 is dictated by the design specifications.

When coil 2 travels within the magnetic field of a permanent magnet (MFPM), an induced electric current arises within coil 2. The intensity of this current is influenced by the velocity of coil 2, as well as its internal electrical resistance and the normal component of the magnetic induction vector (NCMIV) of MFPM perpendicular to the plane of coil 2. This induced current in moving coil 2 subsequently produces an alternating magnetic field around coil 1. The electromotive force of coil 1 resulting from the magnetic flux of the alternating magnetic field created by the induced current in moving coil 2 acts as a signal for carrying useful information about the system parameters.

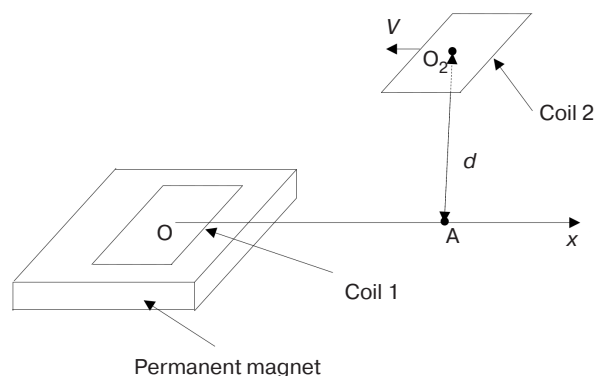


Fig. 1. Schematic of the magnetic system

Various analytical [3–11] and numerical [12–24] techniques have been established for analyzing and calculating the properties of magnetic fields within diverse magnetic systems. The present study employs the equivalent solenoid method to precisely calculate the NCMIV distribution within the MFPM in the plane parallel to the PM upper edge. These findings are then used to create the NCMIV analytical model, which can be used to develop analytical models for magnetic fluxes responsible for generating induced electric currents in conducting coils 1 and 2. The findings regarding the MFPM distribution, as detailed in [11], are expanded upon by additional calculations focusing on the MFPM distribution in the geometric area near permanent magnet.

PROBLEM STATEMENT

Figure 2 depicts the magnetic field produced by a permanent magnet that is uniformly magnetized along the Oz axis. The origin of the coordinate system $Oxyz$ is placed at the magnet base center. Based on Fig. 2, we calculate the normalized MFPM NCMIV, $B_{zn}(x, y, z)$, at point $M(x, y, z)$ located in the specified plane $z \geq H$.

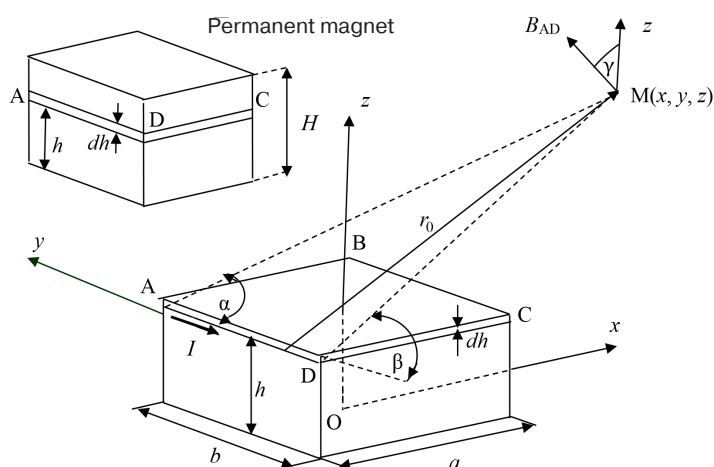


Fig. 2. Geometric layout for calculating MFPM NCMIV

Utilizing the equivalent solenoid approach [3–6], we substitute the permanent magnet with a system of surface electric currents, characterized by a linear density λ , flowing in planes perpendicular to the Oz axis. Then, a surface electric current of strength $I = \lambda dh$ flows along a strip spanning the PM side surface, with a width of dh and situated at

a height h above its base. The strip perimeter is defined by the segments AD, DC, CB, and BA, where $|AD| = |CB| = b$, $|DC| = |BA| = a$.

The magnetic field strength, $B_{AD}(x, y, z)$, created by current I flowing through segment AD at point $M(x, y, z)$ is determined by the following equation [3]:

$$B_{AD}(x, y, z) = \frac{\mu_0}{4\pi} I \frac{\cos \alpha - \cos \beta}{r_0}, \quad (1)$$

where r_0 is the distance from point $M(x, y, z)$ to segment AD.

The magnetic induction vector B_{AD} is directed at an angle γ to the Oz axis.

Based on Fig. 2, the following relationships are derived:

$$-\cos \beta = \frac{y + \frac{b}{2}}{\sqrt{\left(x + \frac{a}{2}\right)^2 + \left(y + \frac{b}{2}\right)^2 + (z - h)^2}}, \quad (2)$$

$$\cos \alpha = \frac{\frac{b}{2} - y}{\sqrt{\left(x + \frac{a}{2}\right)^2 + \left(y - \frac{b}{2}\right)^2 + (z - h)^2}}, \quad (3)$$

$$\cos \gamma = \frac{x + \frac{a}{2}}{\sqrt{\left(x + \frac{a}{2}\right)^2 + (z - h)^2}}. \quad (4)$$

The NCMIV B_{ADz} generated by the electric current flowing through segment AD is determined by the following equation:

$$B_{ADz}(x, y, z) = B_{AD}(x, y, z) \cos \gamma. \quad (5)$$

The NCMIV arising from the electric currents within segments DC, CB, and BA are determined in the same way.

The NCMIV produced by the electric current I flowing through a strip of width dh at a height h above PM is the sum of NCMIVs generated by the currents within segments AD, DC, CB, and BA:

$$B_{hz}(x, y, z) = B_{ADz}(x, y, z) + B_{DCz}(x, y, z) + B_{CBz}(x, y, z) + B_{BAz}(x, y, z). \quad (6)$$

In the following, the AD, DC, CB, and BA symbols denote the PB side surfaces containing corresponding segments.

Taking into account Eqs. (1)–(6), the formula for calculating the normalized NCMIV of the electric current on the side surface AD is derived as follows:

$$\begin{aligned} \frac{4\pi}{\mu_0 \lambda} \int_0^H B_{ADz}(x, y, z) dh = & \left(\arcsin \frac{\alpha_{11} z}{\sqrt{(\alpha_{11}^2 + \beta_1^2)(\beta_1^2 + z^2)}} - \arcsin \frac{\alpha_{11}(z - H)}{\sqrt{(\alpha_{11}^2 + \beta_1^2)(\beta_1^2 + (z - H)^2)}} \right) + \\ & + \left(\arcsin \frac{\alpha_{12} z}{\sqrt{(\alpha_{12}^2 + \beta_1^2)(\beta_1^2 + z^2)}} - \arcsin \frac{\alpha_{12}(z - H)}{\sqrt{(\alpha_{12}^2 + \beta_1^2)(\beta_1^2 + (z - H)^2)}} \right), \end{aligned} \quad (7)$$

where $\alpha_{11} = \frac{b}{2} - y$, $\alpha_{12} = \frac{b}{2} + y$, $\beta_1 = x + \frac{a}{2}$.

The normalized MFPM NCMIV, $B_{zn}(x, y, z)$, equals the sum of the normalized NCMIVs of all electric currents on the PM side surfaces. Taking into account (7), we obtain the following:

$$B_{zn}(x, y, z) = \sum_{i=1}^4 \left(\arcsin \frac{\alpha_{i1}z}{\sqrt{(\alpha_{i1}^2 + \beta_i^2)(\beta_i^2 + z^2)}} - \arcsin \frac{\alpha_{i1}(z-H)}{\sqrt{(\alpha_{i1}^2 + \beta_i^2)(\beta_i^2 + (z-H)^2)}} \right) + \left(\arcsin \frac{\alpha_{i2}z}{\sqrt{(\alpha_{i2}^2 + \beta_i^2)(\beta_i^2 + z^2)}} - \arcsin \frac{\alpha_{i2}(z-H)}{\sqrt{(\alpha_{i2}^2 + \beta_i^2)(\beta_i^2 + (z-H)^2)}} \right) \text{sign}(\beta_i), \quad (8)$$

where $\alpha_{11} = \frac{b}{2} - y$, $\alpha_{12} = \frac{b}{2} + y$, $\beta_1 = x + \frac{a}{2}$ (for AD); $\alpha_{21} = x + \frac{a}{2}$, $\alpha_{22} = \frac{a}{2} - x$, $\beta_2 = y + \frac{b}{2}$ (for DC); $\alpha_{31} = y + \frac{b}{2}$, $\alpha_{32} = \frac{b}{2} - y$, $\beta_3 = \frac{a}{2} - x$ (for CB); $\alpha_{41} = \frac{a}{2} - x$, $\alpha_{42} = \frac{a}{2} + x$, $\beta_4 = \frac{b}{2} - y$ (for BA).

The NCMIV of the electric current on the CB side surface changes sign at $x > \frac{a}{2}$, similar to the NCMIV of the electric current on the BA side surface at $y > \frac{b}{2}$. The change of NCMIV signs is taken into account by the function.

Since the function $B_{zn}(x, y, z)$ is even with respect to both x and y coordinates, determining its distribution within a specified plane only requires evaluating it in one quarter as limited by the inequalities $0 \leq x$ and $0 \leq y$. To simplify $B_{zn}(x, y, z)$ for further investigations, formula (8) is decomposed into a Taylor series. The analysis indicates that formula (8) can be replaced by the more concise and computationally efficient NCMIV analytical model as follows:

$$B_{zn,anal}(x, y, z) = \begin{cases} B_{zn}(0, 0, z) \left[\left(\frac{y}{y_{\max}} \right)^2 - 1 \right]^2 \left[\left(\frac{x}{x_{\max}} \right)^2 - 1 \right]^2 & \text{at } |x| \leq x_{\max}, |y| \leq y_{\max}, \\ 0 & \text{at } |x| > x_{\max}, |y| > y_{\max}, \end{cases} \quad (9)$$

where $x_{\max}, y_{\max}$ are the “dimensions” of the magnetic field in the plane $z = \text{const}$, i.e., the coordinates $x_{\max}, y_{\max}$, which depend on the coordinate z , are chosen so that the MFPM NCMIV becomes equal to zero in the region $|x| \geq x_{\max}, |y| \geq y_{\max}$ of the plane $z = \text{const}$.

In Fig. 3, the graphs of the dependencies of the normalized MFPM NCMIV, $B_{zn}(x, y, z)$, in the plane $z = 13$ mm are presented as a function of the x -coordinate at various y -coordinate values, with the following PM parameters: $a = b = 16$ mm, $H = 8$ mm. The values of $B_{zn}(x, y, z)$ are calculated using the exact formula (8) for graphs 1, 2, and 3 and formula (9) for graphs 4, 5, and 6.

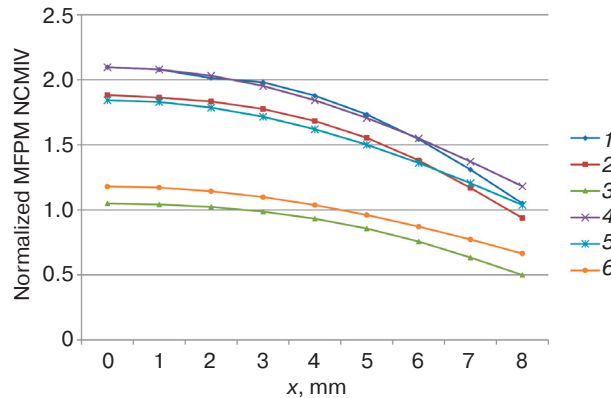


Fig. 3. Graphs of the normalized MFPM NCMIV: 1 is $B_{zn}(x, 0, 13)$, 2 is $B_{zn}(x, 4, 13)$, 3 is $B_{zn}(x, 8, 13)$, 4 is $B_{zn,anal}(x, 0, 13)$, 5 is $B_{zn,anal}(x, 4, 13)$, and 6 is $B_{zn,anal}(x, 8, 13)$

The analysis reveals a significant difference between the calculation results using the exact formula (8) and those using the approximation formula (9) for points in the plane with coordinates $x \geq 6$ mm, $y \geq 6$ mm. The relative error, δ , of the MFPM NCMIV calculations based on the analytical model belongs to the segment $15\% \leq \delta \leq 33\%$.

For points on the plane with coordinates $x \leq 6$ mm, $y \leq 6$ mm, the relative error δ is less than 15%. Moreover, for points on the plane limited by the inequality $r \leq 4$ mm, where r is the distance between the coordinate system origin and the point on the plane, the relative error is $\delta \leq 3.8\%$; for points of the plane limited by the inequality $r \leq 2$ mm, the relative error is $\delta \leq 2.6\%$.

The parameters of the useful signal of the measuring system, including its amplitude and time-dependent shape over its duration, are determined by the rate of change of magnetic fluxes within coils 1 and 2 when the center of coil 2 is above the center of PM cross-section. Since the magnetic fluxes and the rate of their changes are primarily determined by the central part of the MFPM NCMIV as represented by the points in the plane with coordinates limited by the inequalities $x \leq 6$ mm and $y \leq 6$ mm, the error in calculating the useful signal parameters in this case is minimal.

At high values of the MFPM NCMIV coordinates, the MFPM peripheral part decreases rapidly, while the rate of change of the magnetic fluxes determined by the MFPM peripheral part is less than that determined by the MFPM central part.

Error estimations derived from NCMIV calculations employing the approximation formula (9) reveal the potential of the analytical NCMIV model for creating an analytical model for the useful signal of the measuring system. This analytical model achieves a calculation error for useful signal values not exceeding 15%.

CONCLUSIONS

Based on the results of the research, the following can be deduced:

1. An analytical formula for calculating the NCMIV of the parallelepiped-shaped MFPM, which is oriented perpendicular to a specified plane aligned with the PM end face, has been derived by applying the equivalent solenoid technique.
2. The obtained expression for the NCMIV analytical model enables straightforward calculations of the magnetic fluxes within coils 1 and 2.
3. The feasibility of applying the NCMIV analytical model to theoretically develop an analytical model for the useful signal of the measuring system with inductive signal transmission is confirmed.

REFERENCES

1. Zakatov M.M. Theoretical foundations of an information system with inductive signal transmission. *Radiopromyshlennost' = Radio Industry*. 2001;1:29–39 (in Russ.).
2. Zakatov M.M. *Passive Automotive Tire Air Pressure and Temperature Meter*: RF Pat. 2509656. Publ. 20.03.2014 (in Russ.).
3. Sivukhin D.V. *Obshchii kurs fiziki: Elektrichestvo (General Course of Physics: in 5 v. V. 3. Electricity)*. Moscow: Fizmatlit; 2004. 656 p. (in Russ.).
4. Kustov M.S., Druina D.V., Mikhailova O.O., Polyakov I.G. Magnetic field calculation for uniformly magnetized bodies by the methods of equivalent solenoid and magnetic charges. *Vestnik TvGU. Seriya: Fizika = Bulletin of TvGU. Series: Physics*. 2010;8:17–35 (in Russ.).
5. Erlichson H. The magnetic field of a circular turn. *Am. J. Phys.* 1989;57(7):607–608. <https://doi.org/10.1119/1.15955>
6. Vorontsov A.A., Slesarev Yu.N., Karpushin E.V. Mathematical modeling of magnetic fields of two-axis magnetostrictive inclinometers containing constant magnet shaped as rectangular parallelepiped. *Vestnik TGTU = Transactions of the TSTU*. 2013;19(1):25–29 (in Russ.).
7. Vorontsov A.A., Slesarev Yu.N., Karpushin E.V. Mathematical modeling of magnetic fields of two-coordinate magnetostrictive tiltmeters containing a ring or solid permanent magnet. *Izvestiya Penzenskogo gosudarstvennogo pedagogicheskogo universiteta im. V.G. Belinskogo*. 2012;30:467–472 (in Russ.).
8. Ravaut R., Lemarquand G., Depollier C., et al. Analytical calculation of the magnetic field created by permanent-magnet rings. *IEEE Transact. Magnet.* 2008;44(8):1982–1989. <https://doi.org/10.1109/TMAG.2008.923096>
9. Vagin D.V., Gerasimenko T.N., Polyakov P.A. Exact analytical expression for magnetic field induction of rectangular shape sample. *Moscow Univ. Phys.* 2008;63(6):422–422. <https://doi.org/10.3103/S0027134908060131> [Original Russian Text: Vagin D.V., Gerasimenko T.N., Polyakov P.A. Exact analytical expression for magnetic field induction of rectangular shape sample. *Vestnik Moskovskogo universiteta. Seriya 3: Fizika. Astronomiya*. 2008;6:53–55 (in Russ.).]
10. Xiao-Fan G., Yong Y., Xiao-Jing Z. Analytic expression of magnetic field distribution of rectangular permanent magnets. *Appl. Math. Mech.* 2004;25:297–306. <https://doi.org/10.1007/BF02437333>
11. Artamonov A.N., Erzhenin N.R., Zakatov M.M. Analytical model of the normal component of the magnetic induction vector. In: *Proceedings of the 27th International Scientific-Practical Conference "Prevention. Salvation. Help."* Khimki; 2017. P. 5–9 (in Russ.).
12. Rezinkina M.M. Parameters of thin electromagnetic shields that provide a decrease in magnetic induction. *Tech. Phys.* 2014;59(2):155–161. <https://doi.org/10.1134/S1063784214020194> [Original Russian Text: Rezinkina M.M. Parameters of thin electromagnetic shields that provide a decrease in magnetic induction. *Zhurnal tekhnicheskoi fiziki*. 2014;84(2):1–7 (in Russ.).]

13. Rezinkina M.M. Numerical calculation of the magnetic field and magnetic moment of ferromagnetic bodies of complex spatial configuration. *Tech. Phys.* 2009;54(8):1092–1101. <https://doi.org/10.1134/S1063784209080027>
[Original Russian Text: Rezinkina M.M. Numerical calculation of the magnetic field and magnetic moment of ferromagnetic bodies of complex spatial configuration. *Zhurnal tekhnicheskoi fiziki.* 2009;79(8):8–17 (in Russ.).]
14. Slesarev Yu.N., Martens-Atyusheva K.Yu., Rodionov S.V. Research of mathematical models of magnetic fields of two-coordinate magnetostriiction tiltmeters. *Modeli, sistemy, seti v ekonomike, tekhnike, prirode i obshchestve = Models, Systems, Networks in Economics, Technology, Nature and Society.* 2016;2(18):290–298 (in Russ.).
15. Bulzhev E.M., Menshov E.N., Dzhavakhiya G.A. Modeling the field permanent magnet. *Izvestiya Samarskogo nauchnogo tsentra RAN = Izvestiya of Samara Scientific Center of the Russian Academy of Sciences.* 2011;13(4):106–110 (in Russ.).
16. Cherkasova O.A. Research of the magnetic field of the permanent magnet by means of computer modeling. *Geteromagnitnaya mikroelektronika = Heteromagnetic Microelectronics.* 2014;17:112–120 (in Russ.).
17. Matyuk V.F., Osipov A.A., Strelyukhin A.V. Magnetization distribution along the cylindrical rod in longitudinal constant uniform magnetic field. *Russ. Electr. Eng.* 2009;80(8):450–458. <https://doi.org/10.3103/S1068371209080082>
[Original Russian Text: Matyuk V.F., Osipov A.A., Strelyukhin A.V. Magnetization distribution along the cylindrical rod in longitudinal constant uniform magnetic field. *Elektrotehnika.* 2009;8:37–46 (in Russ.).]
18. Andreeva E.G., Tatevosyan A.A., Semina I.A. Research of the axisymmetric model of an open-type magnetic system. *Omskii nauchnyi vestnik = Omsk Scientific Bulletin.* 2010;1(87):110–113 (in Russ.).
19. Lagutin A.S., Grigoriev G.Yu. Magnetic system of the xenon core polarizer. *Prikladnaya fizika = Appl. Phys.* 2023;3:92–95 (in Russ.).
20. Alievskii B.L., Orlov V.L. *Raschet parametrov magnitnykh polei osesimmetrichnykh katushek (Calculation of Magnetic Field Parameters of Axisymmetric Coils).* Moscow: Energoatomizdat; 1983. 112 p. (in Russ.).
21. Camacho J.M., Sosa V. Alternative method to calculate the magnetic field of permanent magnets with azimuthal symmetry. *Revista Mexicana de Física E.* 2013;59(1):8–17. <https://doi.org/10.31349/RevMexFisE.21.020701>
22. Semykina I.Yu., Zav'yalov V.M., Dubkov E.A., Uglova M.B. Analysis of the magnetic field induction distribution in the vicinity of a rectangular coil profile of a wireless power transmission system. In: *Proceedings of the International Scientific and Practical Conference "Advanced Technologies and Materials."* Sevastopol. 2022. P. 235–239 (in Russ.).
23. Li B., Zhang J., Zhao X., Liu B., Dong H. Research on air gap magnetic field characteristics of trapezoidal Halbach permanent magnet linear synchronous motor based on improved equivalent surface current. *Energies.* 2023;16(2):793–805. <https://doi.org/10.3390/en16020793>
24. Osipov V.V., Orlov A.N., Lisenkov V.V., et al. Distribution of the Magnetic Field in the Gap between Two Permanent Magnets: Calculated and Experimental Data and Their Applications. *Instrum. Exp. Tech.* 2023;66(6):995–1002. <https://doi.org/10.1134/S0020441223060064>
[Original Russian Text: Osipov V.V., Orlov A.N., Lisenkov V.V., Maksimov R.N., Shitov V.A. Distribution of the Magnetic Field in the Gap between Two Permanent Magnets: Calculated and Experimental Data and Their Applications. *Pribory i tekhnika eksperimenta.* 2023;6:111–118 (in Russ.). <https://doi.org/10.31857/S003281622306006X>]

СПИСОК ЛИТЕРАТУРЫ

1. Закатов М.М. Теоретические основы информационной системы с индукционной передачей сигнала. *Радиопромышленность.* 2001;1:29–39.
2. Закатов М.М. Пассивный измеритель давления и температуры воздуха в шине колеса автомобиля: пат. 2509656 РФ. Заявка № 2012140219/11; заявл. 20.09.2012, опубл. 20.03.2014. Бюл. 8.
3. Сивухин Д.В. *Общий курс физики: в 5 т. Т. 3. Электричество.* М.: Физматлит; 2004. 656 с.
4. Кустов М.С., Друина Д.В., Михайлова О.О., Поляков И.Г. Расчет магнитных полей однородно намагниченных тел методами эквивалентного соленоида и магнитных зарядов. *Вестник ТвГУ. Серия «Физика».* 2010;8:17–35.
5. Erlichson H. The magnetic field of a circular turn. *Am. J. Phys.* 1989;57(7):607–608. <https://doi.org/10.1119/1.15955>
6. Воронцов А.А., Слесарев Ю.Н., Карпунин Э.В. Математическое моделирование магнитных полей двухкоординатных магнитострикционных наклономеров, содержащих постоянный магнит в форме прямоугольного параллелепипеда. *Вестник Тамбовского ГТУ.* 2013;19(1):25–29.
7. Воронцов А.А., Слесарев Ю.Н., Карпунин Э.В. Математическое моделирование магнитных полей двухкоординатных магнитострикционных наклономеров, содержащих кольцевой или сплошной постоянный магнит. *Известия Пензенского государственного педагогического университета им. В.Г. Белинского.* 2012;30:467–472.
8. Ravaud R., Lemarquand G., Depollier C., et al. Analytical calculation of the magnetic field created by permanent-magnet rings. *IEEE Transact. Magnet.* 2008;44(8):1982–1989. <https://doi.org/10.1109/TMAG.2008.923096>
9. Вагин Д.В., Герасименко Т.Н., Поляков П.А. Точное аналитическое выражение для индукции магнитного поля образца прямоугольной формы. *Вестник Московского университета. Серия 3: Физика. Астрономия.* 2008;6:53–55.
10. Xiao-Fan G., Yong Y., Xiao-Jing Z. Analytic expression of magnetic field distribution of rectangular permanent magnets. *Appl. Math. Mech.* 2004;25:297–306. <https://doi.org/10.1007/BF02437333>
11. Артамонов А. Н., Ерженин Н.Р., Закатов М.М. Аналитическая модель нормальной составляющей вектора магнитной индукции. В сб.: *Материалы XXVII Международной научно-практической конференции «Предупреждение, спасение, помощь».* Химки; 2017. С. 5–9.

12. Резинкина М.М. Выбор параметров тонких электромагнитных экранов для снижения уровней магнитной индукции. *Журн. техн. физики*. 2014;84(2):1–7.
13. Резинкина М.М. Численный расчет магнитного поля и магнитного момента ферромагнитных тел сложной пространственной конфигурации. *Журн. техн. физики*. 2009;79(8):8–17.
14. Слесарев Ю.Н., Мартенс-Атюшева К.Ю., Родионов С.В. Исследование математических моделей магнитных полей двухкоординатных магнитоотражающих наклонеров. *Модели, системы, сети в экономике, технике, природе и обществе*. 2016;2(18):290–298.
15. Булыжев Е.М., Меньшов Е.Н., Джавахия Г.А. Моделирование поля постоянного магнита *Известия Самарского научного центра РАН*. 2011;13(4):106–110.
16. Черкасова О.А. Исследование магнитного поля постоянного магнита с помощью компьютерного моделирования. *Гетеромагнитная микроэлектроника*. 2014;17:112–120.
17. Матюк В.Ф., Осипов А.А., Стрелюхин А.В. Распределение намагниченности вдоль цилиндрического стержня, находящегося в продольном постоянном однородном магнитном поле. *Электротехника*. 2009;8:37–46.
18. Андреева Е.Г., Татевосян А.А., Семина И.А. Исследование осесимметричной модели магнитной системы открытого типа. *Омский научный вестник*. 2010;1(87):110–113.
19. Лагутин А.С., Григорьев Г.Ю. Магнитная система поляризатора ядер ксенона. *Прикладная физика*. 2023;3:92–95.
20. Алиевский Б.Л., Орлов В.Л. *Расчет параметров магнитных полей осесимметричных катушек*. М.: Энергоатомиздат; 1983. 112 с.
21. Camacho J.M., Sosa V. Alternative method to calculate the magnetic field of permanent magnets with azimuthal symmetry. *Revista Mexicana de Física E*. 2013;59(1):8–17. <https://doi.org/10.31349/RevMexFisE.21.020701>
22. Семькина И.Ю., Завьялов В.М., Дубков Е.А., Углова М.Б. Анализ распределения индукции магнитного поля в окрестностях катушки прямоугольного профиля системы беспроводной передачи энергии. В сб.: *Перспективные технологии и материалы: Материалы Международной научно-практической конференции*. Севастополь. 2022. С. 235–239.
23. Li B., Zhang J., Zhao X., Liu B., Dong H. Research on air gap magnetic field characteristics of trapezoidal Halbach permanent magnet linear synchronous motor based on improved equivalent surface current method. *Energies*. 2023;16(2):793–805. <https://doi.org/10.3390/en16020793>
24. Осипов В.В., Орлов А.Н., Лисенков В.В., Максимов Р.Н., Шитов В.А. Распределение магнитного поля в зазоре между двумя постоянными магнитами: расчетные и экспериментальные данные и их приложения. *Приборы и техника эксперимента*. 2023;6:111–118. <https://doi.org/10.31857/S003281622306006X>

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